ABSTRACT: The actual silicon crisis in photovoltaic industry forces the whole value chain to realize more installed Wp photovoltaic power out of every kg Silicon. This is approached through thinner wafers by the wafer producers, higher cell efficiencies by the cell manufacturers, and efficiency enhancements via improved encapsulation schemes at module production. Recently, some alternative encapsulation materials came in consideration, e.g. EVA replacements and cover glass with anti-reflective coatings. Cell technologies are changing as well. Some innovations are already adopted in standard products, but need to be tuned to each other. While the number of parameters involved has increased, it became cumbersome to minimize the optical losses experimentally. Modelling and variation of those parameters is performed to understand their interdependence and to propose optimal parameters sets, which can be used as good starting parameters for experimental probing.

Keywords: anti-reflective coating, modules, optical losses

INTRODUCTION

In this paper a simple algorithm for modelling the optical losses of encapsulated solar cells is developed. The anti-reflective coating (ARC) of the cell, i.e. silicon nitride (SiN_x) single layer, as the actual standard, plays a key role in the encapsulation behaviour. The optical losses in the module after encapsulation are highlighted in detail by calculating the reflectance and absorption of the different layers with the given optical constants of silicon [1] silicon nitride [2], ethylene vinyl acetate EVA [3] and solar glass [4]. Finally, the model is used to demonstrate the limits of the state-of-the art encapsulation technology on next generation silicon cells with improved spectral response (SR).

THEORETICAL APPROACH

EQE of the encapsulated cell

The short circuit current of the encapsulated cell in the module is calculated from the measured external quantum efficiency (EQE) of the cell in air. Together with the cell's reflectance spectrum (either measured or calculated from \( n_{\text{SiN}} \) and \( d_{\text{SiN}} \)) and the calculated losses due to the absorption and reflectance of the module encapsulation, see Fig. 1, the EQE of the module is:

\[
\begin{align*}
EQE_{\text{module}} &= EQE_{\text{cell/air}} \left( \frac{1 - A_{\text{SiN}}} {1 - A_{\text{SiN}}} \right) \ldots \\
&= \frac{(1 - R_{\text{glass/air}})(1 - R_{\text{cell/EVA}})(1 - R_{\text{cell/EVA}})(1 + (1 - A_{\text{encapsulation}})R_{\text{cell/EVA}})}{(1 - R_{\text{cell/EVA}})}.
\end{align*}
\]

The last factor in the numerator is describing the effect of multiple reflections, where higher orders with more than two reflections hitting the surface of the cell are neglected. The reflection at the EVA/glass interface is omitted for the multiple reflections as well, because \( R_{\text{glass/EVA}} \) is below 0.1% for the sensitive range of silicon cells between 300 nm to 1200 nm.

\[
I_{\text{AM1.5}} \quad \Sigma = \text{Encapsulation}
\]

\[
\begin{array}{c}
\text{ARC glass, ca. 110 nm} \\
\text{Cover glass, 3.2 mm} \\
\text{EVA, 0.4 mm} \\
\text{SiN_H ARC, 70 nm} \\
\text{Silicon, 220 µm}
\end{array}
\]

**Fig. 1. Reflectance and absorption losses after encapsulation of a solar cell.**

\( A_{\text{SiN}} \) and \( A_{\text{SiN'}} \) serve to differentiate between alternative silicon nitride layers with different parameters. The reflectance and absorption of the different layers are calculated in detail in the next section. The short circuit photocurrent of the encapsulated cell in the module calculates to:

\[
I_{\text{module}} = \int_{0}^{1200} \frac {q_{\text{e}}} {hc} \cdot EQE_{\text{module}} \cdot I_{\text{AM1.5}} \cdot d\lambda
\]

with \( I_{\text{AM1.5}}(\lambda) \) as the global solar spectrum in steps of 10 nm taken from IEC 60904-3 Ed. 2 [5].

Reflectance

The reflectance of the SiN_c coated cell \( R_{\text{cell/air}} \) in air or \( R_{\text{cell/EVA}} \) in EVA is given by [6]:

\[
R_{\text{cell/air}} = \frac {\hat{n}_1 m_1 + \hat{n}_2 m_2 - m_2 - \hat{n}_1 m_1} {\hat{n}_1 m_1 + \hat{n}_2 m_2 + m_1 + \hat{n}_2 m_2}
\]

with the transfer matrix

\[
\begin{bmatrix}
\hat{n}_1 m_1 + \hat{n}_2 m_2 - m_2 - \hat{n}_1 m_1 \\
\hat{n}_1 m_1 + \hat{n}_2 m_2 + m_1 + \hat{n}_2 m_2
\end{bmatrix}
\]
\[
\begin{aligned}
\left( \frac{m_1}{m_2}, \frac{m_1}{m_3} \right) &= \left[ \cos \delta_{\text{air}} + i \sin \delta_{\text{air}}/n_{\text{air}} \right] \\
\left( \frac{m_1}{m_2}, \frac{m_1}{m_3} \right) &= \left[ i n_{\text{air}} \sin \delta_{\text{air}} + \cos \delta_{\text{air}} \right]
\end{aligned}
\] (3b)

where \( n_i \) are the complex refractive indices for the incident medium \( n_{\text{air}} \), i.e., air or EVA, respectively, for silicon \( n_s \) and for the silicon nitride layer \( n_{\text{Si}} \) in the complex form:

\[
\begin{aligned}
n_i &= n - i k_i \\
\delta &= \frac{2\pi d}{\lambda} n_{\text{Si}}
\end{aligned}
\] (3c)

The above formulas were implemented in an Excel worksheet (enabled to operate with complex arguments by the Add-in “Technical analysis functions”). The reflectance on the thick adjacent layers, e.g. glass/EVA, was calculated by:

\[
\begin{aligned}
R_{\text{glass/EVA}}(\lambda) &= \left[ \left( n_{\text{glass}} - n_{\text{EVA}} \right)^2 + k_{\text{EVA}}^2 \right]^{1/2} \\
&= \frac{n_{\text{glass}} + n_{\text{EVA}}}{n_{\text{glass}} + n_{\text{EVA}}} + k_{\text{EVA}}
\end{aligned}
\] (4)

The reflectance of the encapsulation is not used for the modelling here, but might be the parameter accessible in the experiment. Taking multiple reflections into account, it is:

\[
\begin{aligned}
R_{\text{encaps}} &= R_{\text{glass/air}} + R_{\text{glass/EVA}} + R_{\text{glass/air}} \left( 1 - A_{\text{encaps}} \right) \left( 1 - R_{\text{glass/air}} \right) + \ldots \left( 1 - R_{\text{glass/air}} \right)^n + \ldots \\
&= \frac{R_{\text{glass/air}} \left( 1 - A_{\text{encaps}} \right)}{1 + R_{\text{glass/air}} \left( 1 - A_{\text{encaps}} \right)}
\end{aligned}
\] (5)

**Table:**

<table>
<thead>
<tr>
<th>Material</th>
<th>Reflectance (%)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Flat solar glass (polished Albarino)</td>
<td>12%</td>
</tr>
<tr>
<td>With porous layer ARC (conventional)</td>
<td>20%</td>
</tr>
<tr>
<td>With multilayer ARC (PV-Lite)</td>
<td>15%</td>
</tr>
<tr>
<td>With ARC calculated with d=130nm n=1.23</td>
<td>5%</td>
</tr>
</tbody>
</table>

Fig. 2. Effect of different anti-reflective coatings on the reflectance of the glass/air interface.

from measured glass reflectance and transmittance spectra provided by [4] and using Eq. 5 to separate the effect of the two different interfaces.

Recently, Wohlgemuth et al. [7] showed an increase of 2.4% in the power output at STC and an additional increase of 1.8% in the performance ratio using cover glasses with anti-reflective coating. The glass suppliers prospect a power increase of 2.7% to 3.1% at standard test conditions (STC). We calculated an possible increase of 3.7% for a single layer ARC on a polished glass surface with a refractive index of \( n = 1.23 \) and a thickness of 130 nm. As a design rule, one should optimise the encapsulated cell to uncoated glass first and then adjust the ARC to it, because the effects from the glass ARC are much smaller then those from the ARC of the cell. The multi-layer ARC in Fig. 2 might be an exception, because of its high reflectance beyond the anti-reflective window.

**Absorption**

The absorption \( A_{\text{SiN}} \) of the Si\(_{\text{N}}\), is given by

\[
1 - \exp(-\alpha_{\text{SiN}} d_{\text{SiN}}) = 4\pi k_{\text{SiN}}/\lambda.
\]

The absorption in EVA and glass is calculated using thicknesses of 0.4 mm and 3.2 mm, respectively, which are well established values in industrial practice regarding module reliability. The optical constants for EVA and glass are taken from [3],[4].

\[
1 - A_{\text{encaps}} = \exp(-\alpha_{\text{EVA}} d_{\text{EVA}}) \exp(-\alpha_{\text{glass}} d_{\text{glass}})
\]

The module’s “zero-depth concentrator”-effect resulting from the backscattering of the incident light at the white back sheet material (Tedlar\textsuperscript{®}) from the spacings between the cells in the module were determined experimentally to 0.4% additional \( I_{\text{sc}} \) current per 1 mm spacing for 156 mm cells in the range of 3-5 mm cell spacing. This is in good agreement with the calculated value of 1.3% for 3 mm spacings as provided by McIntosh et al. [8].

The reflections of the soldered interconnectors covering typically 3% of the cell area add only 0.13%. Both effects sum up to 1.4% in total for 3 mm 156 mm square cells.

**Texturization**

The effect of the texturization of the cell’s surface is approximated by assuming a second impact of the light reflected at the textured surface, i.e. \( R_{\text{cell/air}} \) is replaced by \( R_{\text{cell/air}} \).

This approximation is already a good guess to describe the difference in the experimental EQE between the nearly flat, alkaline etched, multi crystalline cell in Fig. 3 and the alkaline textured mono crystalline cell as measured by Fraunhofer ISE Freiburg. The shading loss from the metal front grid is 8% for both cells.

An “improved cell” was modelled in Fig. 3 by setting the quantum efficiency to be constant in the wavelength range starting from 550 nm to sketch the effect of an ideal emitter, what increases the I\textsubscript{sc} by +1.6%. In the infrared region the collection efficiency was fixed for values of \( \lambda > 700 \) nm and the decrease of the silicon absorption was replaced by the effect of an Lambertian back reflector on the back of the cell, which can be described with a path length enhancement of \( 4\pi f_{\text{R}} \) or +10% in \( I_{\text{sc}} \) in the present case or a total of 40 mA/cm\(^2\).
Fig. 3. Measured and calculated EQE as used in the calculation of the encapsulation effects

RESULTS

SiN$_x$ layer optimization

Fig. 4 shows the variation of the short circuit current for an encapsulated multi cell (Q6L) as a function of the refractive index and thickness of the SiN$_x$ layer.

![Graph showing EQE variation with SiN$_x$ layer optimization](image)

The optimum value of $n_{SiNx}=2.22$ at 632.8 nm and $d=67$ nm in Fig. 4 confirm the results found by Doshi [10] and Ekai [11]. An independent set of cells with a total $I_{sc}$ variation of 6% as a result from differences in the wafer quality, i.e. variations in the infrared region of the spectral response (SR) resulted in a deviation of ±0.01 for the found optimal refractive index and ±1nm for the found thickness. While Doshi found a refractive index of $n=2.23$ and $d=68$ nm for the encapsulated cell, Ekai et al. [11] calculated $n=2.20$ and $d=67$ nm for flat cells and $n=2.10$ and $d=70$ nm for V-grooved cell surfaces, respectively.

Taking the zero-depth concentrator effects into account (1.4% enhancement), the encapsulation gain for weakly or non-textured multi cells is 4.8% when using an ideal ARC glass and it is 1.1% without ARC glass.

Encapsulation factor

Fig. 5 shows the principal agreement between calculated and experimental data.

![Graph showing encapsulation factor comparison](image)

A detailed analysis of the optical losses is plotted in Fig. 6 for a multi cell. The contribution of each loss mechanism was calculated by setting all other losses to zero. In consequence the total loss is slightly smaller than the direct sum of the single losses.

![Graph showing optical losses and encapsulation factor](image)

The cell’s reflectance is the major loss either in air as in the module. The encapsulated cell is gaining mainly in the range between 400 and 600 nm, where the index matching between silicon and EVA by the SiN$_x$ layer results in reflection reductions of 50%.
Encapsulation of textured cells

Fig. 7 shows the difference in the cell reflectance in air and with EVA for the non-textured multi cell (0% textured area), an intermediate case with 50% textured area and the 100% textured cell (e.g. mono cells).

Fig. 7. Degree of texturization and its effect on the optical losses and the encapsulation factor

Fig. 7 demonstrates the decrease of the encapsulation gain with increasing degree of texturization. It decreases in steps of 2%.

After optimisation of the SiN parameters for each degree of texturization, see curve descriptions in Fig. 7, the encapsulation gain decreases again, see Table 1, but the Isc will increase overall, i.e. results in a better utilization of the given wafer material.

<table>
<thead>
<tr>
<th>texturization</th>
<th>real cell</th>
<th>improved cell (model)</th>
</tr>
</thead>
<tbody>
<tr>
<td>0% textured</td>
<td>31.1 mA/cm²</td>
<td>37.8 mA/cm²</td>
</tr>
<tr>
<td>SiNx parameters</td>
<td>n=2.22, d=67 nm</td>
<td>n=2.16, d=68 nm</td>
</tr>
<tr>
<td>SiNx absorption in air</td>
<td>-1.3%</td>
<td>-1.3%</td>
</tr>
<tr>
<td>SiNx absorption in module</td>
<td>-1.3%</td>
<td>-0.2%</td>
</tr>
<tr>
<td>EVA absorption</td>
<td>-1.3%</td>
<td>-2.4%</td>
</tr>
<tr>
<td>glass absorption</td>
<td>-0.8%</td>
<td>-0.9%</td>
</tr>
<tr>
<td>ARC/glass/air reflection</td>
<td>-0.7%</td>
<td>-0.8%</td>
</tr>
<tr>
<td>cell reflection</td>
<td>-9.6%</td>
<td>-8.9%</td>
</tr>
<tr>
<td>total</td>
<td>-10.5%</td>
<td>-10.9%</td>
</tr>
<tr>
<td>encapsulation gain</td>
<td>3.4%</td>
<td>2.1%</td>
</tr>
</tbody>
</table>

Table 1: Isc losses in air and in the module with a standard cell and the improved cell from Fig. 3 without considering the zero-depth concentrator effect

For the fully textured cells, the optimum for the SiNx antireflective layer shifted to a lower refractive index of n=2.04 and a slightly lower thickness of d=66 nm, because it is more effective to reduce the SiNx absorption through lowering n and d, while reduction of the cell reflectance is nearly obsolete for textured cells.

UV absorption in EVA has become the major loss mechanism for the improved cell in Table 1. The effect of the ideal emitter, i.e. 1.6% increase in the UV, is completely blocked by the EVA.

CONCLUSION

A simple model was developed to describe the complex system of an encapsulated solar cell, which can be implemented in spread-sheet program. The optimum parameters for the SiNx layer of non-textured cells are $n_{SiNx}=2.22$ and $d_{SiNx}=67$ nm. The zero-depth concentrator effect adds 1.4% in Isc for 3 mm spaced 156 mm square cells on white Tedlar®. The total encapsulation gain sums up to 1.1%. With increasing degree of texturization the optimum refractive index of the SiN decreases. For fully textured cells (e.g. mono cells) it is $n_{SiNx}=2.04$. ARC cover glass does not need adjustments on the cell side and an ideal gain of 3.7% was calculated. The effect of shallower emitters is likely to be limited by the EVA absorption. Implications on the optical losses for non-STC conditions are considered in an accompanying paper [12].

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REFERENCES

[8] K. R. McIntosh et al. Prog. in Photovoltaics 14, 2006, p.167